

# Analysis of Breakdown Characteristics of Field-Plate AlGaIn/GaN HEMTs with a High- $k$ Passivation Layer

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## ABSTRACT

It is well known that the introduction of field plate increases the breakdown voltage  $V_{br}$  of AlGaIn/GaN HEMTs. As another way to improve  $V_{br}$ , using a high- $k$  passivation layer is proposed. So, in this study, we combine the two factors and analyzed the breakdown characteristics of field-plate AlGaIn/GaN HEMTs as parameters of its length  $L_{FP}$  and the relative permittivity of passivation layer  $\epsilon_r$ . It is shown that the enhancement of  $V_{br}$  with increasing  $\epsilon_r$  is more significant when  $L_{FP}$  is relatively short. There is an optimum value of  $L_{FP}$  to obtain the highest  $V_{br}$ , and it is around 0.2 and 0.3  $\mu\text{m}$  when the gate-to-drain distance is 1.5  $\mu\text{m}$ . When  $L_{FP} = 0.3 \mu\text{m}$  and  $\epsilon_r$  takes a high value of 50, the electric field between the field-plate edge and the drain becomes rather uniform, and  $V_{br}$  becomes about 400 V, which corresponds to an effective electric field of 2.7 MV/cm between gate and drain.

**Keywords:** GaN HEMT, breakdown voltage, field plate, high- $k$  passivation layer

## 1 INTRODUCTION

AlGaIn/GaN HEMTs are now attractive for applications to high-power microwave devices and high-power switching devices [1, 2]. It is well known that introducing a field plate improves the power performance of AlGaIn/GaN HEMTs as well as GaAs-based devices [3-5]. This is because the field plate reduce current collapse [6, 7], and also increase the breakdown voltage  $V_{br}$  [8-10].

To increase  $V_{br}$ , the introduction of passivation layer with high permittivity (high- $k$  layer) is also considered [11-13]. In a previous work [12, 13], we analyzed off-state breakdown characteristics in AlGaIn/GaN HEMTs as a parameter of relative permittivity of passivation layer  $\epsilon_r$ , and found that  $V_{br}$  increased as  $\epsilon_r$  increased.

Therefore, in this study, we combine the two structures and analyzed the breakdown characteristics of field-plate AlGaIn/GaN HEMTs with a high- $k$  passivation layer, and investigate how  $V_{br}$  is enhanced.

## 2 PHYSICAL MODEL

Fig.1 shows a device structure analyzed in this study. The gate length  $L_G$  is 0.3  $\mu\text{m}$  and the gate-to-drain distance

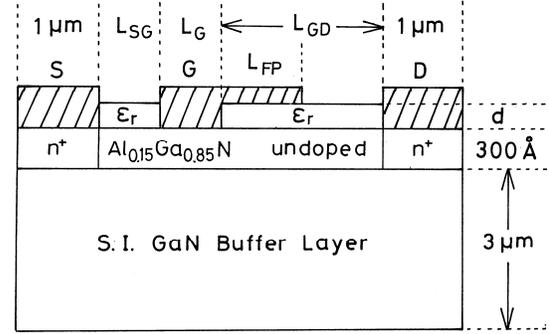


Figure 1: Device structure analyzed in this study.

$L_{GD}$  is 1.5  $\mu\text{m}$ . The thickness of passivation layer  $d$  is 0.1  $\mu\text{m}$ . The field-plate length  $L_{FP}$  is varied between 0 and 1  $\mu\text{m}$ . The relative permittivity of the passivation layer  $\epsilon_r$  is varied between 1 and 50. In a semi-insulating buffer layer, we consider a shallow donor, a deep donor, and a deep acceptor [14-16]. As an energy level of the deep acceptor, we consider  $E_C - 2.85 \text{ eV}$  ( $E_V + 0.6 \text{ eV}$ ). For impurity compensation, we consider  $E_C - 0.5 \text{ eV}$  as an energy level of the deep donor. The deep-acceptor density  $N_{DA}$  is set rather high of  $10^{17} \text{ cm}^{-3}$ .

Basic equations to be solved are Poisson's equation including ionized deep-level terms and continuity equations for electrons and holes including a carrier generation rate by impact ionization and carrier loss rates via the deep levels [13, 17-19]. These equations are expressed as follows.

1) Poisson's equation

$$\nabla \cdot (\epsilon \nabla \psi) = -q(p - n + N_{DI} + N_{DD}^+ - N_{DA}^-) \quad (1)$$

2) Continuity equations for electrons and holes

$$\nabla \cdot \mathbf{J}_n = -qG + q(R_{DD} + R_{DA}) \quad (2)$$

$$\nabla \cdot \mathbf{J}_p = qG - q(R_{DD} + R_{DA}) \quad (3)$$

where  $N_{DI}^+$  and  $N_{DA}^-$  are the ionized deep-donor and deep-acceptor densities, respectively.  $R_{DD}$  and  $R_{DA}$  represent carrier recombination rates via the deep donors and the deep acceptors, respectively.  $G$  is a carrier generation rate by impact ionization, and given by

$$G = (\alpha_n |J_n| + \alpha_p |J_p|) / q \quad (4)$$

where  $\alpha_n$  and  $\alpha_p$  are ionization rates for electrons and holes, respectively, and expressed as

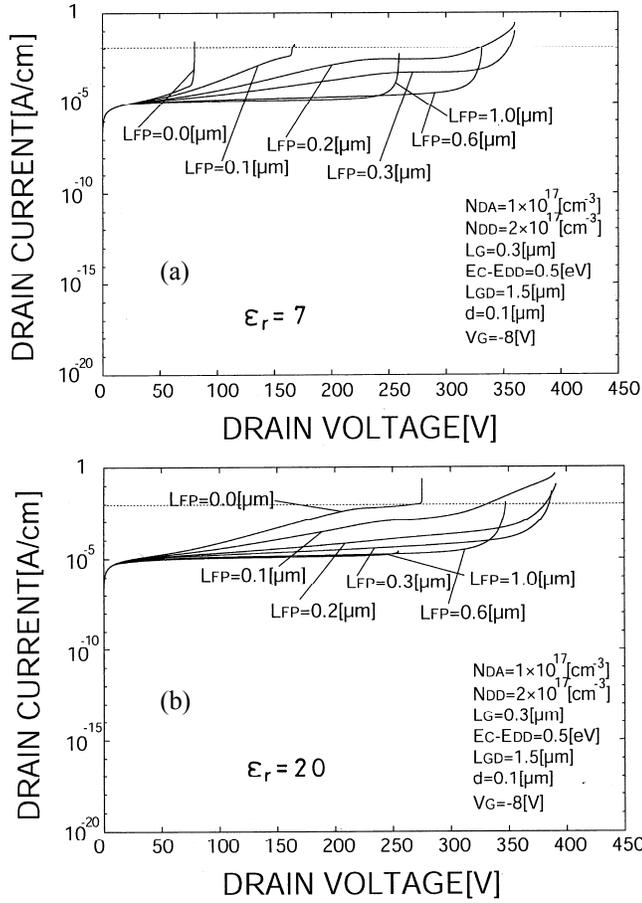


Figure 2: Calculated  $I_D$ - $V_D$  curves as a parameter of  $L_{FP}$ . (a)  $\epsilon_r = 7$ , (b)  $\epsilon_r = 20$ . The dotted lines show 1 mA/mm.

$$\alpha_n = A_n \exp(-B_n / |E|) \quad (5)$$

$$\alpha_p = A_p \exp(-B_p / |E|) \quad (6)$$

where  $E$  is the electric field.  $A_n$ ,  $B_n$ ,  $A_p$ , and  $B_p$  are deduced from [20]. The above basic equations are put into discrete forms and solved numerically.

### 3 RESULTS AND DISCUSSIONS

Figs.2(a) and (b) show a comparison of calculated drain current  $I_D$  – drain voltage  $V_D$  curves of AlGaIn/GaN HEMTs as a parameter of  $L_{FP}$  between the two cases with  $\epsilon_r = 7$  (SiN) and 20, respectively. The gate voltage is  $-8$  V, which corresponds to an off state. In both cases, the drain current usually increases suddenly, showing breakdown (except for  $\epsilon_r = 20$  and  $L_{FP} = 0.1 \text{ } \mu\text{m}$ ). This is due to impact ionization of carriers. Overall, the breakdown voltage seems to be higher in the case of  $\epsilon_r = 20$ . Particularly, it is higher at relatively short  $L_{FP}$ .

Fig.3 shows a comparison of electric field profiles at the heterojunction interface for  $L_{FP} = 0$  between the two cases with  $\epsilon_r = 7$  and 20. In the case of  $\epsilon_r = 7$ , the increase in the

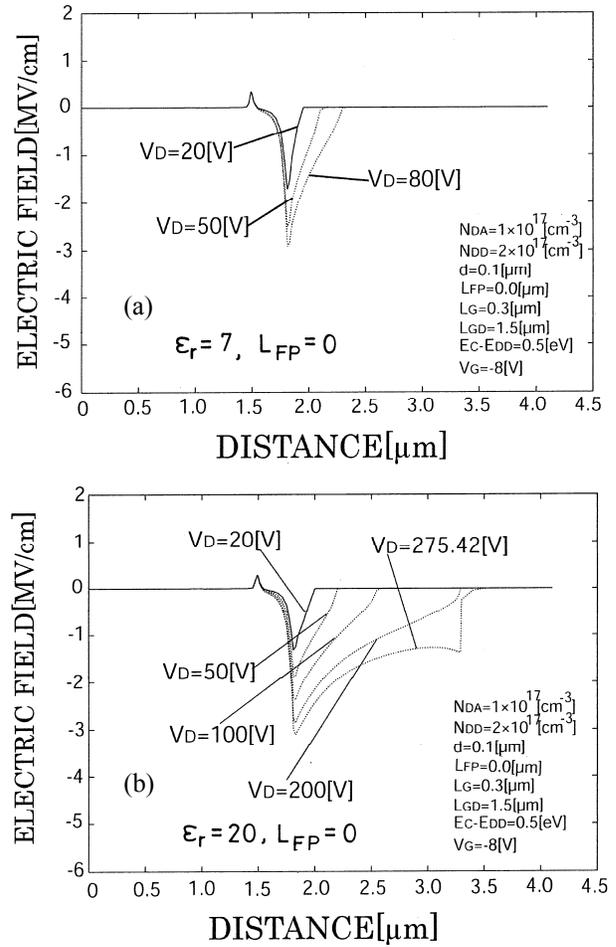


Figure 3: Electric field profiles along the heterojunction interface.  $L_{FP} = 0$ . (a)  $\epsilon_r = 7$ , (b)  $\epsilon_r = 20$ .

drain voltage is entirely applied along the drain edge of the gate, leading to the breakdown at about 80 V. On the other hand, in the case of  $\epsilon_r = 20$ , the electric field at the drain edge of the gate is reduced, and the high field region extends to the drain. Hence, the breakdown voltage increases to about 250 V. Here the breakdown voltage is defined as a drain voltage where the drain current becomes 1 mA/mm.

Fig.4 shows a comparison of electric field profiles at the heterojunction interface for  $L_{FP} = 0.1 \text{ } \mu\text{m}$  between the two cases with  $\epsilon_r = 7$  and 20. In the case of  $\epsilon_r = 7$ , the reduction of electric field at the drain edge is not so significant, leading to the abrupt increase in the drain current at  $V_D = 168$  V. On the other hand, in the case of  $\epsilon_r = 20$ , the electric field at the drain edge of gate is greatly reduced, and the electric field between the gate and drain remains around 3 MV/cm at  $V_D = 390$  V, which is the breakdown voltage. In this case, as shown in Fig.2(b), the drain current does not show an abrupt increase, but increase gradually to reach 1 mA/mm. Here, the buffer leakage current determines the breakdown voltage.

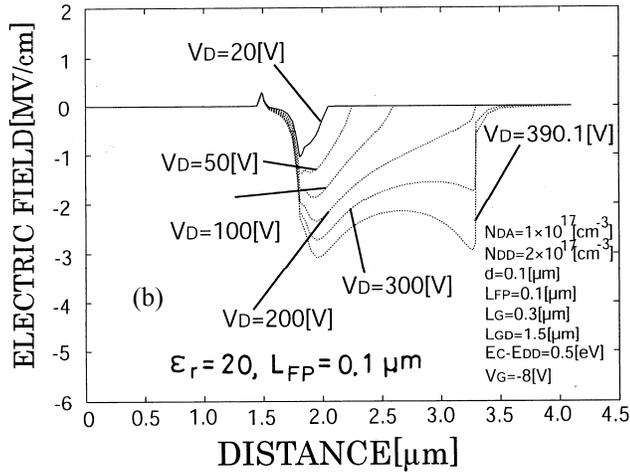
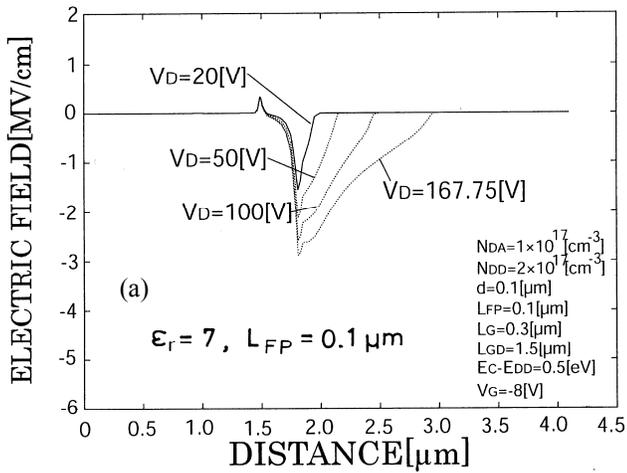


Figure 4: Electric field profiles along the heterojunction interface.  $L_{FP} = 0.1 \mu\text{m}$ . (a)  $\epsilon_r = 7$ , (b)  $\epsilon_r = 20$ .

Fig.5 shows the breakdown voltage  $V_{br}$  as a function of the field-plate length  $L_{FP}$ , with  $\epsilon_r$  as a parameter. It is seen that  $V_{br}$  is higher when  $\epsilon_r$  is higher, particularly in the region where  $L_{FP}$  is relatively short.  $V_{br}$  becomes low when  $L_{FP}$  becomes relatively long. This is because the distance between the field-plate edge and the drain becomes short and the electric field in this region becomes very high. Hence, there is an optimum field-plate length to obtain high  $V_{br}$ , and it is around 0.2 and 0.3  $\mu\text{m}$  here.

Fig.6 shows a comparison of electric field profiles at  $L_{FP} = 0.3 \mu\text{m}$  between the two cases with  $\epsilon_r = 7$  and 50. In both cases, electric field at the drain edge of the gate is reduced, and the breakdown voltage is determined by the electric field between the field-plate edge and the drain. In the case of  $\epsilon_r = 50$ , the electric field in this region is more uniform. This is due to the high- $k$  dielectric, and hence  $V_{br}$  becomes higher (about 400 V) than that for lower  $\epsilon_r$ .

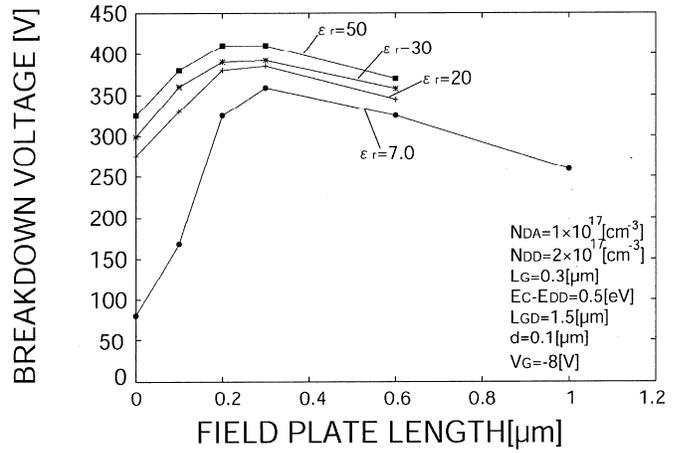


Figure 5: Breakdown voltage  $V_{br}$  versus field-plate length  $L_{FP}$  curves, with  $\epsilon_r$  as a parameter.

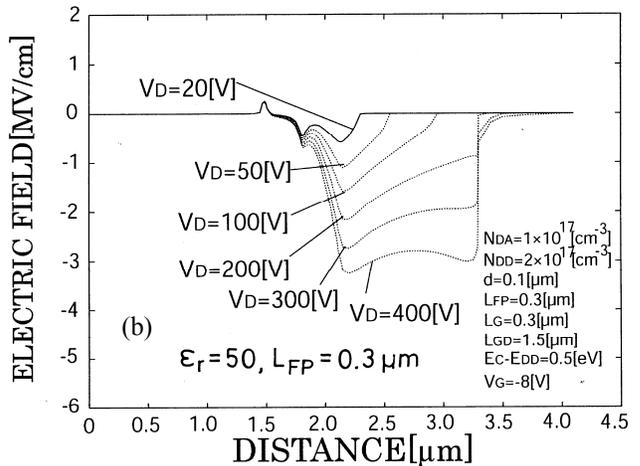
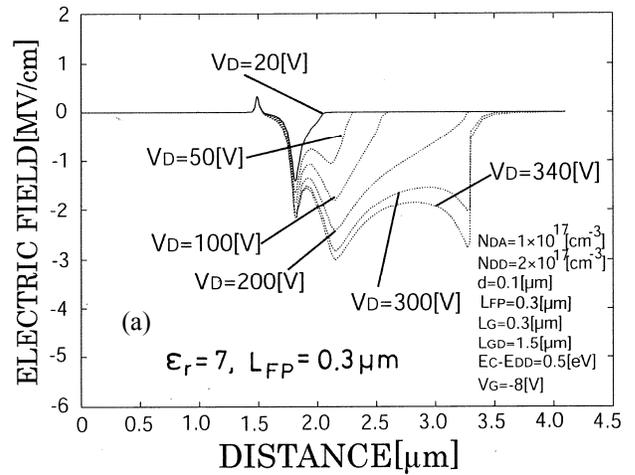


Figure 6: Electric field profiles along the heterojunction interface.  $L_{FP} = 0.3 \mu\text{m}$ . (a)  $\epsilon_r = 7$ , (b)  $\epsilon_r = 50$ .

## 4 CONCLUSION

In summary, we have performed a two-dimensional analysis of breakdown characteristics of field-plate AlGaIn/GaN HEMTs as parameters of its length  $L_{FP}$  and the relative permittivity of passivation layer  $\epsilon_r$ . It has been shown that  $V_{br}$  is higher when  $\epsilon_r$  becomes higher, particularly in the region where  $L_{FP}$  is short. This is because the electric field at the drain edge of gate is effectively reduced even for short  $L_{FP}$ . It has also been shown that there is an optimum value of  $L_{FP}$  to obtain highest  $V_{br}$ , and it is around 0.2 and 0.3  $\mu\text{m}$  when the gate-to-drain distance is 1.5  $\mu\text{m}$ . When  $\epsilon_r$  has a high value of 50, the electric field between the field-plate edge and the drain becomes rather uniform, and  $V_{br}$  reaches about 400 V, which corresponds to an effective electric field of 2.7 MV/cm between the gate and the drain.

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