

# Simulation of Buffer-Related Current Slump in AlGaIn/GaN HEMTs

Kazushige Horio

Faculty of Systems Engineering, Shibaura Institute of Technology  
307 Fukasaku, Minuma-ku, Saitama 337-8570, Japan, horio@sic.shibaura-it.ac.jp

## ABSTRACT

Transient simulations of AlGaIn/GaN HEMTs are performed in which a three-level compensation model is adopted for a semi-insulating buffer-layer, where a shallow donor, a deep donor and a deep acceptor are considered. Quasi-pulsed  $I$ - $V$  curves are derived from the transient characteristics, and are compared with steady-state  $I$ - $V$  curves. It is shown that so-called current slump is more pronounced when the deep-acceptor density in the buffer layer is higher and when an off-state drain voltage is higher, because trapping effects become more significant. It is suggested that to minimize current slump in AlGaIn/GaN HEMTs, an acceptor density in a semi-insulating buffer layer should be made low.

**Keywords:** AlGaIn/GaN HEMT, deep level, current slump, device simulation

## 1 INTRODUCTION

Recently, AlGaIn/GaN HEMTs have received great attention because of their potential applications to high power and high temperature microwave devices [1]. However, slow current transients are often observed even if the drain voltage or the gate voltage is changed abruptly [2]. This is called drain lag or gate lag, and is problematic in circuit applications. The slow transients mean that the dc  $I$ - $V$  curves and the RF  $I$ - $V$  curves become quite different, resulting in lower RF power available than that expected from the dc operation [1],[2]. This is called power slump or current slump. These are serious problems, and there are many experimental works reported on these phenomena [1-6]. But, only a few theoretical works have been reported for GaN-based FETs [6,7], where effects of a donor-type surface state (near the valence band) on gate lag and pulsed  $I$ - $V$  curves of AlGaIn/GaN HEMTs are studied [6], and a bulk deep-acceptor effect ( $\sim 1$  eV above the midgap of GaN) is studied for gate lag in AlGaIn/GaN HEMTs [7]. But, the type of traps and their energy levels seemed to be artificial. In this work, we have made simulations of AlGaIn/GaN HEMTs with a semi-insulating buffer layer in which trap levels based on experiments are considered, as in our previous work on GaN MESFETs [8], and showed that the lag phenomena and the current slump could be reproduced. Additionally, we have studied dependence of current slump on the impurity densities in the buffer layer and on an off-state drain voltage.

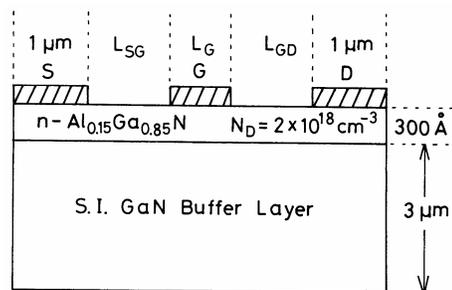


Fig.1. Modeled AlGaIn/GaN HEMT analyzed in this study.

## 2 PHYSICAL MODEL

Fig.1 shows a device structure analyzed in this study. The gate length  $L_G$  is set to  $0.3 \mu\text{m}$ . As a model for the semi-insulating buffer layer, we use a three level compensation model which includes a shallow donor, a deep donor and a deep acceptor. Some experiments show that two levels ( $E_C - 1.7$  eV,  $E_C - 2.85$  eV) are associated with current slump in GaN-based FETs with a semi-insulating buffer layer [2], so that we use energy levels of  $E_C - 2.85$  eV (or  $E_V + 0.6$  eV) for the deep acceptor and of  $E_C - 1.7$  eV for the deep donor. Other experiments show shallower energy levels for the deep donor [9,10], and hence we vary the deep donor's energy level ( $E_{DD}$ ) as a parameter. Here, the deep-donor density ( $N_{DD}$ ) and the deep-acceptor density ( $N_{DA}$ ) are typically set to  $5 \times 10^{16} \text{ cm}^{-3}$  and  $2 \times 10^{16} \text{ cm}^{-3}$ , respectively. The shallow-donor density in the buffer layer  $N_{Di}$  is set to  $10^{15} \text{ cm}^{-3}$ .

Basic equations to be solved are expressed as follows.

1) Poisson's equation

$$\nabla^2 \psi = -\frac{q}{\epsilon} (p - n + N_D + N_{Di} + N_{DD}^+ - N_{DA}^-) \quad (1)$$

2) Continuity equations for electrons and holes

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_n - (R_{n,DD} + R_{n,DA}) \quad (2)$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_p - (R_{p,DD} + R_{p,DA}) \quad (3)$$

where

$$R_{n,DD} = C_{n,DD} N_{DD}^+ n - e_{n,DD} (N_{DD} - N_{DD}^+) \quad (4)$$

$$R_{n,DA} = C_{n,DA} (N_{DA} - N_{DA}^-) n - e_{n,DA} N_{DA}^- \quad (5)$$

$$R_{p,DD} = C_{p,DD} (N_{DD} - N_{DD}^+) p - e_{p,DD} N_{DD}^+ \quad (6)$$

$$R_{p,DA} = C_{p,DA} N_{DA}^- p - e_{p,DA} (N_{DA} - N_{DA}^-) \quad (7)$$

3) Rate equations for the deep levels

$$\frac{\partial}{\partial t} (N_{DD} - N_{DD}^+) = R_{n,DD} - R_{p,DD} \quad (8)$$

$$\frac{\partial}{\partial t} N_{DA}^- = R_{n,DA} - R_{p,DA} \quad (9)$$

where  $N_{DD}^+$  and  $N_{DA}^-$  represent ionized densities of deep donors and deep acceptors, respectively.  $C_n$  and  $C_p$  are the electron and hole capture coefficients of the deep levels, respectively,  $e_n$  and  $e_p$  are the electron and hole emission rates of the deep levels, respectively, and the subscript (DD, DA) represents the corresponding deep level.

The above basic equations are put into discrete forms and are solved numerically. We have calculated the drain-current responses when the drain voltage  $V_D$  and/or the gate voltage  $V_G$  are changed abruptly.

### 3 CURRENT TRANSIENTS

Fig.2 shows calculated drain-current responses when the drain voltage  $V_D$  is raised abruptly from 0 V to 20 V or when  $V_D$  is lowered from 20 V to 10 V, where the gate voltage  $V_G$  is kept constant (0 V). Here, three cases with different  $E_C - E_{DD}$  are shown. When  $V_D$  is raised, the drain currents overshoot steady-state values, because electrons are injected into the buffer layer, and deep traps there need certain time to capture these electrons. On the other hand, when  $V_D$  is lowered, the drain currents remain at low values for some periods and begin to increase slowly, showing drain lag behavior. This is due to the slow response of deep donors. It is understood that the drain currents begin to increase as the deep donors begin to emit electrons, and hence the response is faster for shallower  $E_{DD}$ . In fact, the current rise time is roughly consistent with the deep donor's electron-emission time constant given by  $1/e_{n,DD}$ , which becomes  $3.9 \times 10^{-5}$  s and  $9.8 \times 10^{-3}$  s for  $E_C - E_{DD} = 0.5$  eV and 1.0 eV, respectively. The above overshoot and undershoot behavior is also reported experimentally in AlGaIn/GaN HEMTs [2],[6].

We have next calculated a case when  $V_D$  and  $V_G$  are both changed abruptly. Fig.3 shows calculated turn-on characteristics ( $E_C - E_{DD}$  is 1.0 eV) when  $V_G$  is changed from the threshold voltage  $V_{th}$  to 0 V. The off-state drain voltage  $V_{Doff}$  is 20 V, and the parameter is the on-state drain voltage  $V_{Don}$ . The characteristics are similar to those in

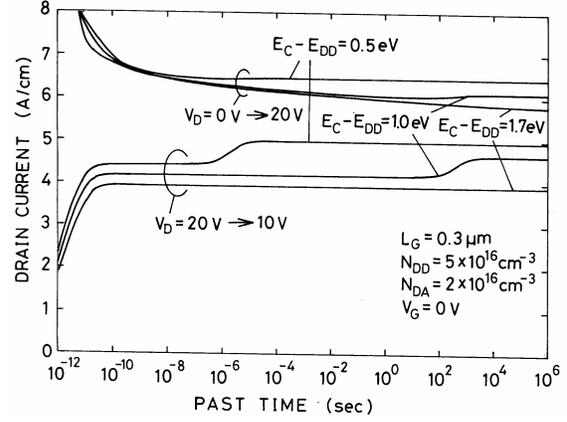


Fig.2. Comparison of drain-current responses of AlGaIn/GaN HEMT as a parameter of deep donor's energy level  $E_{DD}$  when  $V_D$  is raised abruptly from 0 V to 20 V (upper) or when  $V_D$  is lowered abruptly from 20 V to 10 V (lower).  $V_G = 0$  V.  $N_{DD} = 5 \times 10^{16} \text{ cm}^{-3}$  and  $N_{DA} = 2 \times 10^{16} \text{ cm}^{-3}$ .

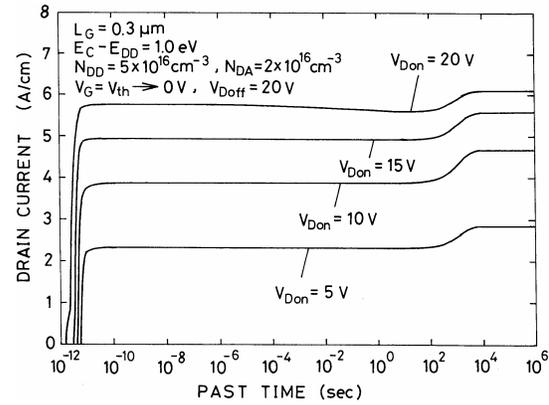


Fig.3. Calculated turn-on characteristics of AlGaIn/GaN HEMT when  $V_G$  is changed from threshold voltage  $V_{th}$  to 0 V, with on-state drain voltage  $V_{Don}$  as a parameter. Off-state drain voltage  $V_{Doff} = 20$  V.  $E_C - E_{DD} = 1.0$  eV.  $N_{DD} = 5 \times 10^{16} \text{ cm}^{-3}$  and  $N_{DA} = 2 \times 10^{16} \text{ cm}^{-3}$ .

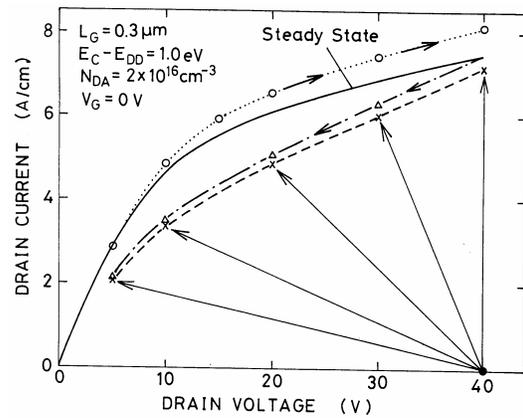


Fig.4. Steady-state  $I$ - $V$  curve ( $V_G = 0$  V; solid line) and quasi-pulsed  $I$ - $V$  curves for AlGaIn/GaN HEMT.  $E_C - E_{DD} = 1.0$  eV. (x):  $V_{Doff} = 40$  V and  $V_{Goff} = V_{th}$  ( $t = 10^{-8}$  s), (o):  $V_D$  is raised from 0 V ( $t = 10^{-9}$  s), ( $\Delta$ ):  $V_D$  is lowered from 40V ( $t = 10^{-8}$  s).

Fig.2, and hence the change of  $V_D$  is regarded as essential in this case. Fig.4 shows calculated  $I_D$ - $V_D$  curves. In this figure, we plot by point (x) the drain current at  $t = 10^{-8}$  s after the gate voltage is switched on. This is obtained from the turn-on characteristics, and this curve corresponds to a quasi-pulsed  $I$ - $V$  curve with pulse width of  $10^{-8}$  s. (We are also plotting other quasi-pulsed  $I$ - $V$  curves when only  $V_D$  is changed, which reflect overshoot and undershoot.) It is seen that the drain currents in the pulsed  $I$ - $V$  curve are rather lower than those in the steady state. This clearly indicates that the current slump could occur due to the slow response of deep levels in the semi-insulating buffer layer. This type of current reduction is commonly observed experimentally in AlGaIn/GaN HEMTs.

## 4 CURRENT SLUMP

### 4.1 Dependence on Deep-Acceptor Density

We have studied dependence of calculated  $I$ - $V$  curves and drain-current responses on the deep-level densities ( $N_{DD}$ ,  $N_{DA}$ ) in the buffer layer. We have found that these characteristics are almost independent of  $N_{DD}$  under a condition that  $N_{DD}$  is higher than  $N_{DA}$  and  $E_C - E_{DD}$  is the same, and that these are mainly determined by  $N_{DA}$ . This is because in this condition, the ionized deep-donor density  $N_{DD}^+$ , which acts as an electron trap, becomes nearly equal to  $N_{DA}$  under equilibrium. Hence, we will show  $N_{DA}$  dependence of the characteristics.

Fig.5 shows calculated  $I_D$ - $V_D$  curves of AlGaIn/GaN HEMTs with  $N_{DA} = 5 \times 10^{15} \text{ cm}^{-3}$  or  $10^{17} \text{ cm}^{-3}$ , where  $N_{DD} = 2 \times 10^{17} \text{ cm}^{-3}$  and  $E_C - E_{DD}$  is 1.0 eV. The solid lines are steady-state  $I$ - $V$  curves. The dashed lines are quasi-pulsed  $I$ - $V$  curves (pulse width of  $10^{-8}$  s) derived from the calculated turn-on characteristics, as mentioned before. It is seen that the steady-state drain currents are higher for lower  $N_{DA}$ . This is because the current via the buffer layer becomes larger for lower  $N_{DA}$  due to less steep barrier at the channel-buffer interface. It is also clearly seen that the current reduction in the pulsed  $I$ - $V$  curves is more pronounced for higher  $N_{DA}$ . This is because, as mentioned before, the ionized deep-donor density  $N_{DD}^+$ , which acts as an electron trap, becomes nearly equal to  $N_{DA}$  under equilibrium, and hence the trapping effect (or the resulting current slump) should become more pronounced for higher  $N_{DA}$ . Here, it should be mentioned that for lower  $N_{DA}$ , the current slump could be weakened, but the threshold voltage shifts toward negative bias because of the higher current density via the buffer layer. Therefore, there may be a trade-off relationship between reducing the current slump and obtaining sharp current cutoff.

### 4.2 Dependence on Off-State Drain Voltage

Next, we describe dependence of current slump on the off-state drain voltage  $V_{Doff}$ . Fig.6 shows calculated steady-state  $I_D$ - $V_D$  curve and quasi-pulsed  $I$ - $V$  curves (with pulse

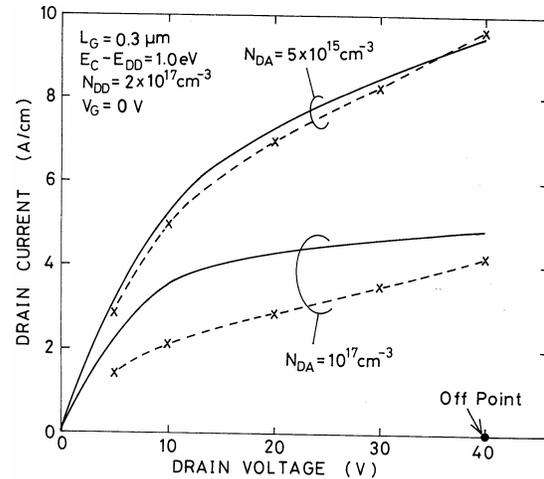


Fig.5. Steady-state  $I$ - $V$  curves ( $V_G = 0$  V; solid lines) and quasi-pulsed  $I$ - $V$  curves (x ;  $t = 10^{-8}$  s) for AlGaIn/GaN HEMTs with different  $N_{DA}$  ( $5 \times 10^{15} \text{ cm}^{-3}$ ,  $10^{17} \text{ cm}^{-3}$ ). Initial point is shown by (●).  $E_C - E_{DD} = 1.0$  eV and  $N_{DD} = 2 \times 10^{17} \text{ cm}^{-3}$ .

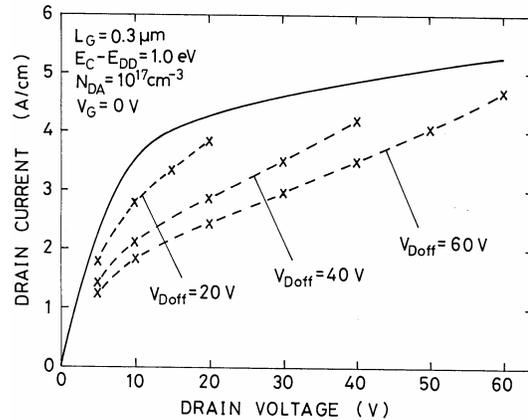


Fig.6. Steady-state  $I$ - $V$  curve ( $V_G = 0$  V; solid line) and quasi-pulsed  $I$ - $V$  curves (x ;  $t = 10^{-8}$  s) for AlGaIn/GaN HEMT, with different off-state drain voltage  $V_{Doff}$ .  $E_C - E_{DD} = 1.0$  eV.  $N_{DD} = 2 \times 10^{17} \text{ cm}^{-3}$  and  $N_{DA} = 10^{17} \text{ cm}^{-3}$ .

width of  $10^{-8}$  s) as a parameter of  $V_{Doff}$ , which are derived from calculated turn-on characteristics as described before. Here,  $N_{DD} = 2 \times 10^{17} \text{ cm}^{-3}$ ,  $N_{DA} = 10^{17} \text{ cm}^{-3}$ , and  $E_C - E_{DD} = 1.0$  eV. It is seen that for higher  $V_{Doff}$ , the drain currents in the pulsed  $I$ - $V$  curves become lower at a given  $V_D$ , indicating that the current slump is more pronounced for higher  $V_{Doff}$ . This is understood as follows. For higher  $V_{Doff}$ , electron densities in the buffer layer become higher particularly under the gate and the gate-to-drain region, because electrons are injected into the buffer layer by the applied drain bias. These electrons are captured by the deep donors, and hence the ionized deep-donor density  $N_{DD}^+$  becomes lower (negative space-charge densities are higher) there for higher  $V_{Doff}$ . Hence, when  $V_G$  is switched on and

$V_D$  is lowered from higher  $V_{Doff}$ , the drain current remains at a lower value, resulting in more pronounced current slump. This tendency is also reported experimentally in AlGaIn/GaN HEMTs [11].

## 5 CONCLUSION

Transient simulations of AlGaIn/GaN HEMTs have been performed in which a three level compensation model is adopted for the semi-insulating buffer layer, where a shallow donor, a deep donor and a deep acceptor are considered. Quasi-pulsed  $I-V$  curves have been derived from the transient characteristics. It has been shown that the lag phenomena and current slump could be reproduced. It has also been shown that the current slump is more pronounced when the deep-acceptor density in the buffer layer is higher and when the off-state drain voltage is higher, because the change of ionized deep-donor density becomes larger and hence the trapping effects become more significant. It is suggested that to minimize current slump in AlGaIn/GaN HEMTs, an acceptor density in a semi-insulating buffer layer should be made low, although there may be a trade-off relationship between reducing current slump and obtaining sharp current cutoff.

## REFERENCES

[1] U. K. Mishra, P. P. Parikh and Y.-F. Wu, "AlGaIn/GaN HEMTs — An overview of device operation and applications", Proc. IEEE, vol.90, pp.1022-1031, 2002.

[2] S. C. Binari, P. B. Klein and T. E. Kazior, "Trapping effects in GaN and SiC Microwave FETs", Proc. IEEE, vol.90, pp.1048-1058, 2002.

[3] M. A. Khan, M. S. Shur, Q. C. Chen and J. N. Kuznia, "Current/voltage characteristics collapse in AlGaIn/GaN heterostructure insulated gate field

effect transistors at high drain bias", Electron Lett., vol.30, pp.2175-2176, 1994.

[4] P. B. Klein, J. A. Freitas, Jr., S. C. Binari and A. E. Wickenden, "Observation of deep traps responsible for current collapse in GaN metal semiconductor field effect transistors", Appl. Phys. Lett., vol.75, pp.4016-4018, 1999.

[5] P. McGovern, J. Benedikt, P. J. Tasker, J. Powell, K. P. Hilton, J. L. Glasper, R. S. Balmer, T. Martin, M. J. Uren, "Analysis of DC-RF dispersion in AlGaIn/GaN HFET's using pulsed  $I-V$  and time-domain waveform measurements", 2005 MTT-S Microwave Symposium Digest, pp.503-506, 2005.

[6] G. Meneghesso et al., "Surface-related drain current dispersion effects in AlGaIn/GaN HEMTs", IEEE Trans. Electron Devices, vol.51, pp.1554-1561, 2004.

[7] N. Braga et al., "Simulation of gate lag and current collapse in GaN heterojunction field effect transistors", Proc. IEEE CSIC Symp., pp.287-290, 2004.

[8] K. Horio, K. Yonemoto, H. Takayanagi and H. Nakano, "Physics-based simulation of buffer-trapping effects on slow current transients and current collapse in GaN field effect transistors", J. Appl. Phys., vol.98, no.12, pp.124502 1-7, 2005.

[9] W. Kruppa, S. C. Binari and K. Doverspike, "Low-frequency dispersion characteristics of GaN HFETs", Electron. Lett., vol.31, pp.1951-1952, 1995.

[10] H. Morkoc, Nitride Semiconductors and Devices, Springer-Verlag, 1999.

[11] A. Koudymov et al., "Dynamic current-voltage characteristics of III-N HFETs", IEEE Electron Device Lett., vol.24, pp.680-682, 2003.